

Anomalies and Entanglement Entropy

Tatsuma Nishioka (University of Tokyo)

based on a work with A. Yarom (Technion)
(to appear)

- An **order parameter** for various phase transitions
 - Confinement/deconfinement (like Polyakov loop) [TN-Takayanagi 06, Klebanov-Kutasov-Murugan 07, ...]
 - Classification of phases [Calabrese-Cardy 04, Kitaev-Preskill 06, Levin-Wen 06, ...]
- Reconstruction of bulk geometry from entanglement
 - Similarity between MERA and AdS space [Swingle 09, Nozaki-Ryu-Takayanagi 12, ...]
 - 1st law of entanglement and linearized Einstein equation of GR [Bhattacharya-Nozaki-Takayanagi-Ugajin 12, Blanco-Casini-Hung-Myers 13, Nozaki-Numasawa-Prudenziati-Takayanagi 13, Lashkari-McDermott-Raamsdonk 13, Faulkner-Guica-Hartman-Myers-Raamsdonk 13, ...]

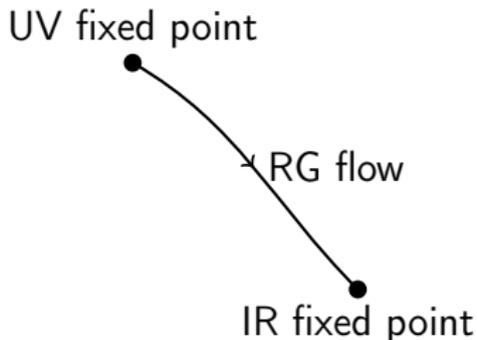
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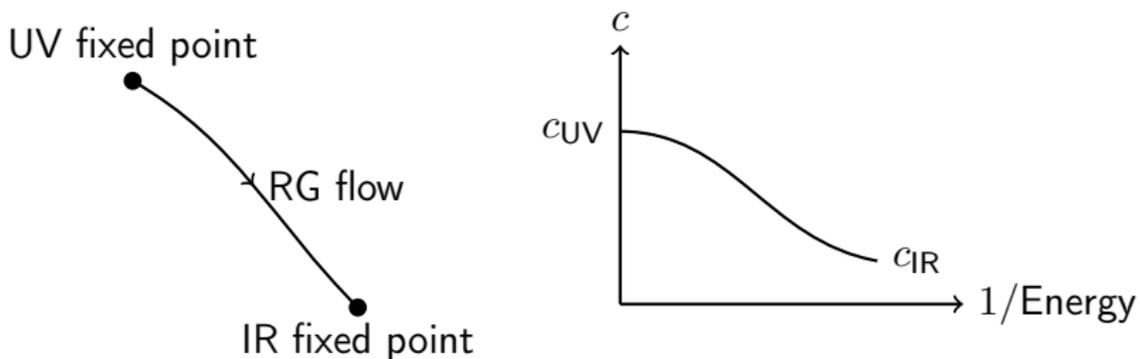
- Entanglement entropy as a measure of degrees of freedom
- Construct a monotonic function $c(\text{Energy})$ of the energy scale
 - Entropic c -theorem in two dimensions [Casini-Huerta 04]
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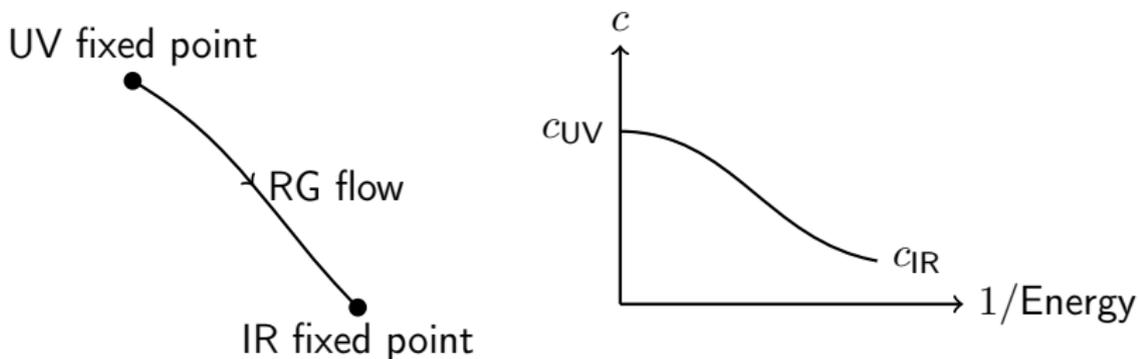
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- Conformal anomalies in even d dimensions: [cf. Myers-Sinha 10]

$$S_A = \frac{c_{d-2}}{\epsilon^{d-2}} + \cdots + c_0 \log \epsilon + \cdots, \quad c_0 \sim \text{central charges}$$

- Relation to sphere partition function $F \equiv (-1)^{\frac{d-1}{2}} \log Z[\mathbb{S}^d]$:
[Casini-Huerta-Myers 11]

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- Gravitational anomaly in CFT_2 with $c_L \neq c_R$

[Wall 11, Castro-Detournay-Iqbal-Perlmutter 14]

(Holography: [Guo-Miao 15, Azeyanagi-Loganayagam-Ng 15])

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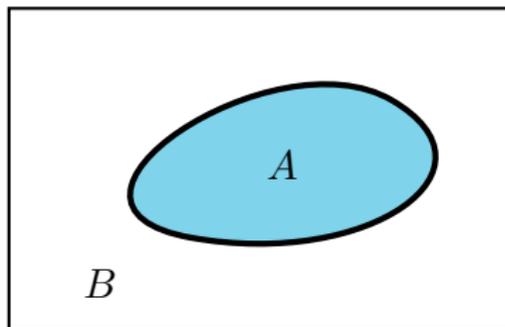
How about chiral and (mixed-)gravitational anomalies in other dimensions?

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- 3 Consistent gravitational anomalies
- 4 Other anomalies

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Definition of entanglement entropy

- Divide a system to A and $B = \bar{A}$: $\mathcal{H}_{tot} = \mathcal{H}_A \otimes \mathcal{H}_B$

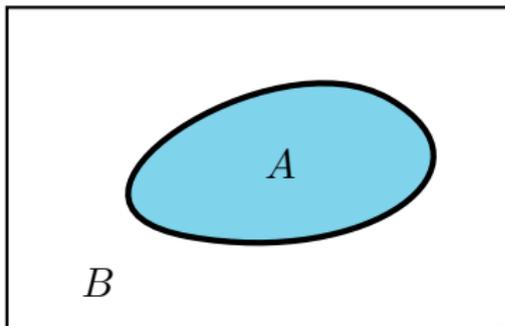


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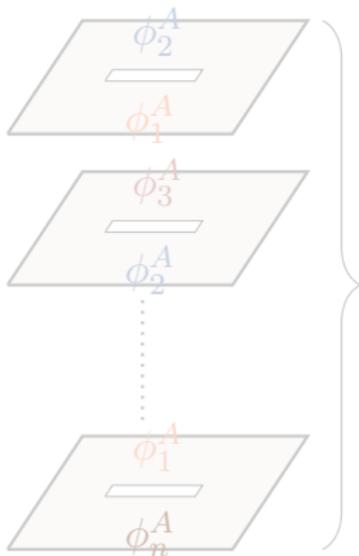
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$$S_A = \lim_{n \rightarrow 1} \frac{1}{1-n} \log \text{tr}_A \rho_A^n$$



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Entanglement entropy

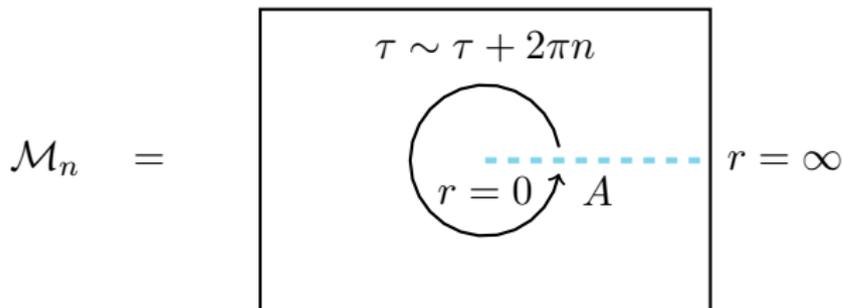
$$S_A = \lim_{n \rightarrow 1} \frac{W_n - nW_1}{n - 1}$$

- $W_n = -\log Z[\mathcal{M}_n]$: the Euclidean partition function
- \mathcal{M}_n : the n -fold cover with a surplus angle $2\pi(n - 1)$ around the entangling surface $\Sigma \equiv \partial A$

The n -fold cover \mathcal{M}_n

- Suppose A is a **semi-infinite line** in two dimensions

$$A = \{0 \leq r < \infty, \tau = 0\}$$

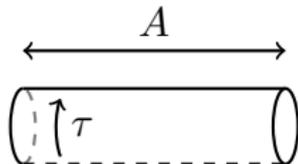


- Use a **partition function W** reproducing flavor and (mixed-)gravitational anomalies
- Evaluate W on the n -fold cover \mathcal{M}_n that is an S^1 fibration over an entangling region A



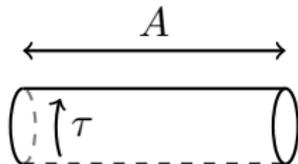
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- Calculate the variation of the entanglement entropy with W

- 1 Entanglement entropy in QFT
- 2 Weyl anomalies
- 3 Consistent gravitational anomalies
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- CFT is classically invariant under the **Weyl rescaling** $\delta_\sigma g_{\mu\nu} = 2\sigma g_{\mu\nu}$ because of $T^\mu{}_\mu = 0$:

$$\delta_\sigma W = \frac{1}{2} \int d^d x \sqrt{g} \sigma T^\mu{}_\mu = 0$$

- The **Weyl anomalies** in even d dimensions:

$$T^\mu{}_\mu = 2(-1)^{d/2} a E - \sum_i c_i I_i$$

E : the Euler density, I_i : Weyl invariants in d dimensions
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$$\delta_\sigma S_A = \lim_{n \rightarrow 1} \frac{\delta_\sigma W_n - n \delta_\sigma W_1}{n - 1}$$

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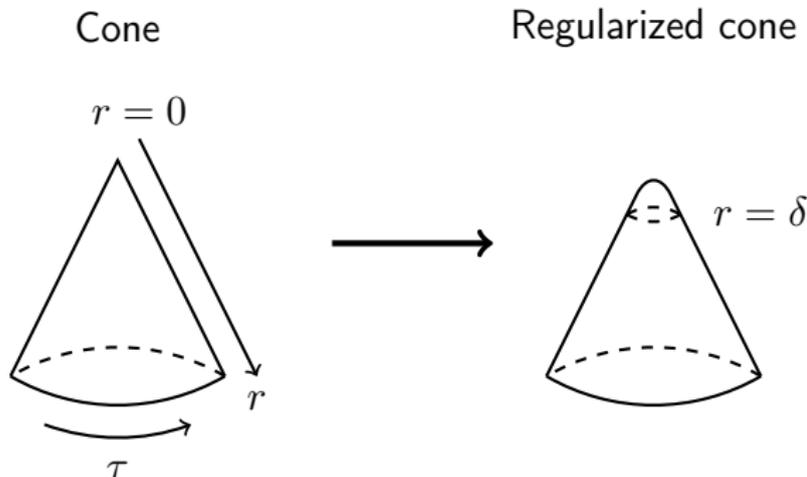
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$$\begin{aligned} \delta_\sigma W_n - n \delta_\sigma W_1 &= \frac{1}{2} \int_{\mathcal{M}_n} \sigma [\langle T^\mu{}_\mu \rangle_{\mathcal{M}_n} - \langle T^\mu{}_\mu \rangle_{\mathcal{M}_1}] \\ &= (n - 1) \sigma \int_\Sigma f(a, c_i, g_{\mu\nu}) + O((n - 1)^2) \end{aligned}$$

because locally $\mathcal{M}_n \simeq \mathcal{M}_1$ away from Σ

- In even dimension:

$$\delta_\sigma S_A = \sigma c_0, \quad c_0 = \int_\Sigma f(a, c_i, g_{\mu\nu})$$

- It gives the **logarithmic divergence**: [Holzhey-Larsen-Wilczek 94, Ryu-Takayanagi 06, Solodukhin 08, Myers-Sinha 10]

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- CFT₂ with central charges $c_L \neq c_R$ has a gravitational anomaly

$$\nabla_\mu T^{\mu\nu} = (c_L - c_R)X^\nu \neq 0$$

- The boost (rotation) of the line A is **not a symmetry**

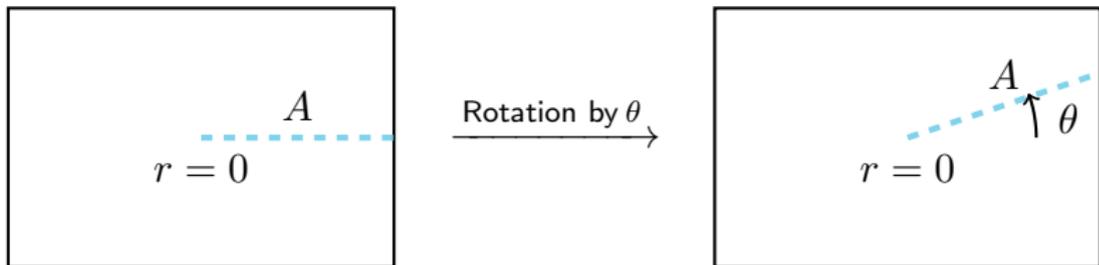


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$$\delta W = \int \Omega_{2k}$$

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 \mathcal{P}_{2k+2}

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I_{2k+1} : Chern-Simons $2k + 1$ -form

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$$W = W_{\text{gauge-inv}} + W_{\text{anom}}$$

- The **consistent anomaly** for W_{anom} can be fixed by the anomaly inflow mechanism [Callan-Harvey 85, Jensen-Loganayagam-Yarom 12]

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$$I_{CS} = c_g \text{tr} \left[\Gamma \wedge d\Gamma + \frac{2}{3} \Gamma \wedge \Gamma \wedge \Gamma \right]$$

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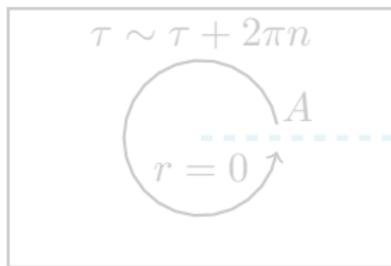
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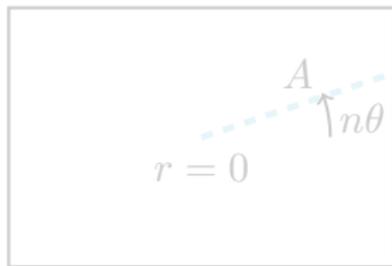
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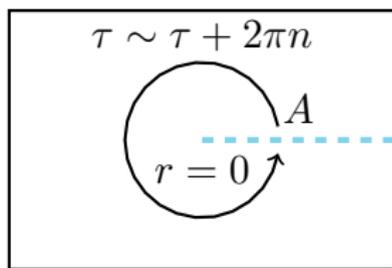
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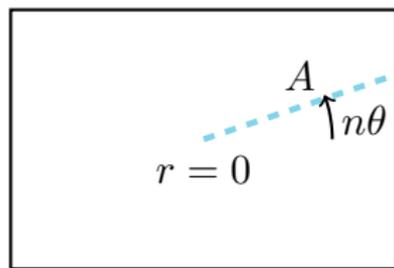
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- Suppose a product space

$$\mathcal{M}_n = \mathcal{C}_n \times \mathcal{N}$$

\mathcal{C}_n : two-dimensional cone, \mathcal{N} : a (compact) two-manifold

- The variation under the diffeomorphism:

$$\delta_\xi W_n = i \int_{\mathcal{C}_n} \alpha c_m \Phi (\partial_\mu \xi^\nu) d\Gamma^\mu{}_\nu$$

with $\Phi \equiv \int_{\mathcal{N}} F$

- Same as the 2d gravitational anomaly with $c_g = \alpha c_m \Phi$

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- The variation under the diffeomorphism:

$$\delta_\xi W_n = i \int_{\mathcal{C}_n} \alpha c_m \Phi (\partial_\mu \xi^\nu) d\Gamma^\mu{}_\nu$$

with $\Phi \equiv \int_{\mathcal{N}} F$

- Same as the 2d gravitational anomaly with $c_g = \alpha c_m \Phi$

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$$\mathcal{P}_{6d} = c_m F^2 \text{tr}(\mathcal{R}^2) + c_a \text{tr}(\mathcal{R}^2)^2 + c_b \text{tr}(\mathcal{R}^4)$$

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Is it always possible to extract coefficients c_m , c_a and c_b from the entanglement entropy?

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Gravitational and mixed anomalies under the boost

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- Anomaly polynomial in higher dimensions:

$$\mathcal{P}_d = \sum_i c_i F^{k_i} \prod_{n=1} \text{tr} (\mathcal{R}^{2n})^{m_n^i}$$

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- 1 Entanglement entropy in QFT
- 2 Weyl anomalies
- 3 Consistent gravitational anomalies
- 4 Other anomalies

- $U(1)$ polygon anomaly in $d = 2m$ dimensions

$$\mathcal{P} = cF^m$$

- For $A \rightarrow A + d\lambda$,

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- Systematic treatment of anomalies in entanglement entropy using the **anomaly inflow mechanism**
- Our method is valid for **any QFT**
- Can be applied as well in higher even dimensions (including **mixed anomalies**), but how about other anomalies such as **parity anomaly** in odd dimensions?