

Ultrasonic non destructive testing advance lecture

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ELyT school

Content

- Element of acoustic
- Few words on piezoelectricity
- US transducer.
- Applications:
 - Medicine
 - High temperature US characterization
 - Reverberation

Element of acoustic and piezoelectricity

ElyT School
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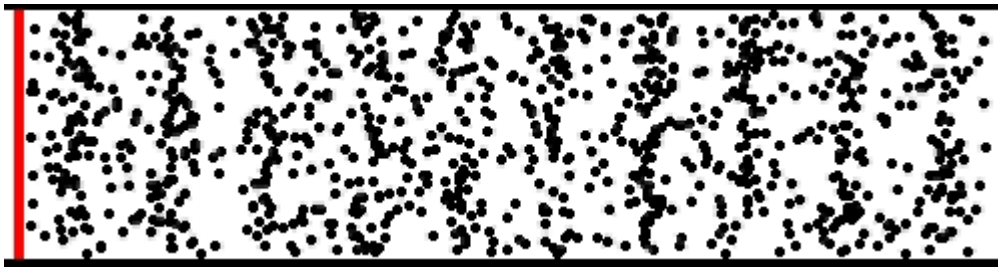
- Outline

- Acoustic waves in fluids
- One dimensional acoustic waves
- Fundamental laws of acoustics
- Energy and acoustics
- Transmission and reflexion

Acoustic waves

The wave vector $\vec{k} = \frac{\omega}{c} \mathbf{u}$

Bulk waves



Longitudinal waves :
Displacement // propagation
(fluids & solids)



Transverse waves
Displacement \perp propagation
(solids only)

Ultrasonic celerity

$$c = \sqrt{\frac{1}{\rho_0 \chi}} \quad \text{In fluids}$$

In solids

$$c_L = \sqrt{\frac{M}{\rho_0}}$$

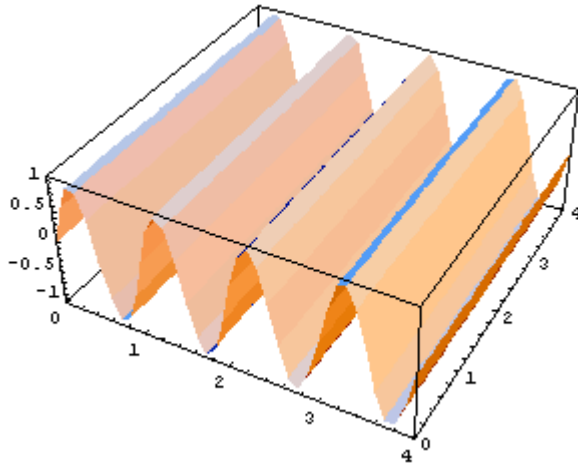
$$c_T = \sqrt{\frac{G}{\rho_0}}$$

$$M = K + \frac{4}{3}G$$

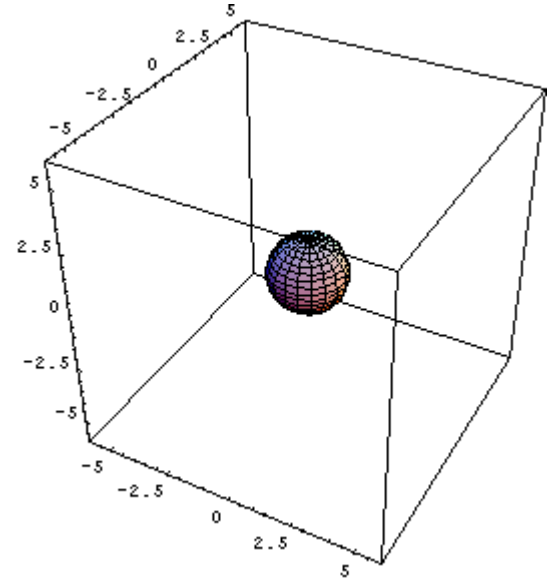
Material	Density [Kg/m ³]	Longitudinal celerity [m/s]	Shear velocity [m/s]
Air (0 degree)	1,293	331	
Air (20 degree)	1,20	344	
Alcohol	790	1207	
Water (pure)	998	1480	
Aluminum	2790	6320	3130
Steel	7800	5900	3200

Waveform

Plane waves



Spherical waves



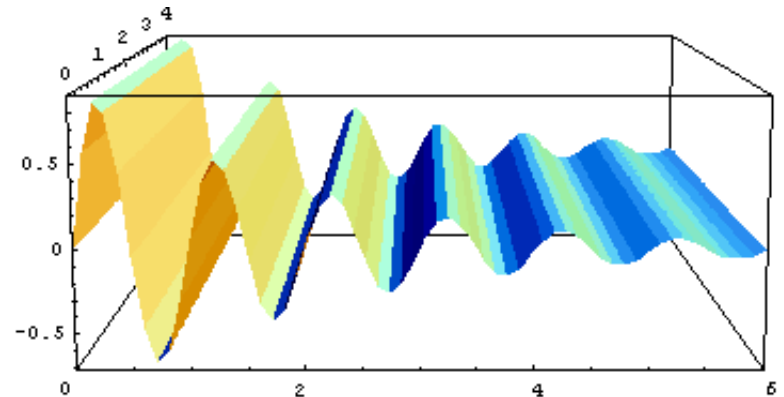
$$u_x = U_M \cos\left(\omega t - \frac{\omega}{c} x\right)$$

Attenuation: diffusion, viscosity

$$u_x = U_M e^{-\alpha x} \cos\left(\omega t - \frac{\omega}{c} x\right)$$

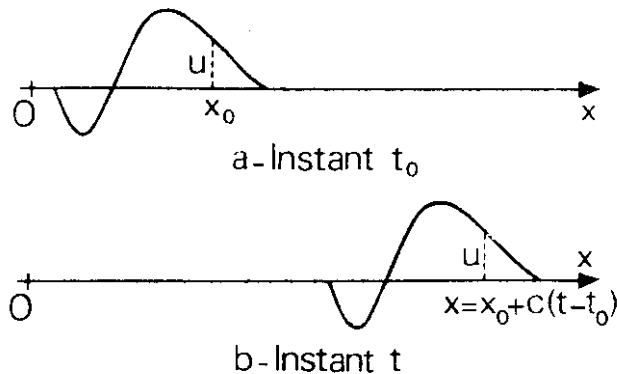
↑
Attenuation factor

$$\alpha = 8.686 \times a \omega^2 \quad \text{dB/m}$$



Water $\alpha = 0.22 \text{ dB/m @ 1MHz}$

Mathematical expression of a plane wave



*Progression d'un ébranlement.
Repéré par le niveau u au point x_0 et à l'instant t_0 , l'ébranlement, se propageant à la vitesse c , atteint, à l'instant t , le point x tel que $x = x_0 + c(t - t_0)$.*

$$x = x_0 + c(t - t_0)$$

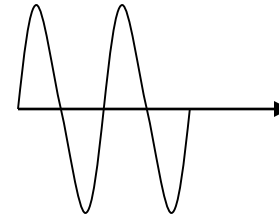
$$x - ct = x_0 - ct_0 \text{ ou } t - \frac{x}{c} = t_0 - \frac{x_0}{c}$$

$$u(x, t) = F\left(t - \frac{x}{c}\right) = f((x - ct)) \text{ Forward propagation } x \nearrow$$

$$u(x, t) = G\left(t + \frac{x}{c}\right) = g((x + ct)) \text{ Backward propagation } x \searrow$$

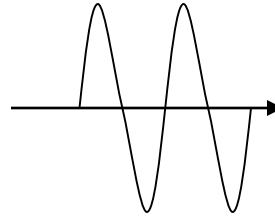
Acoustic wave generated @ $x=0$

$$u_M = A \cos \omega t \quad \text{with} \quad \omega = \frac{2\pi}{T}$$



At a position x

$$u(x, t) = A \cos \omega \left(t - \frac{x}{c} \right)$$



Usefull parameters:

$$\lambda = cT = \frac{c}{f}$$

Wavelength

$$k = \frac{2\pi}{\lambda} = \frac{\omega}{c}$$

Wave number

$$u = A \cos(\omega t - kx)$$

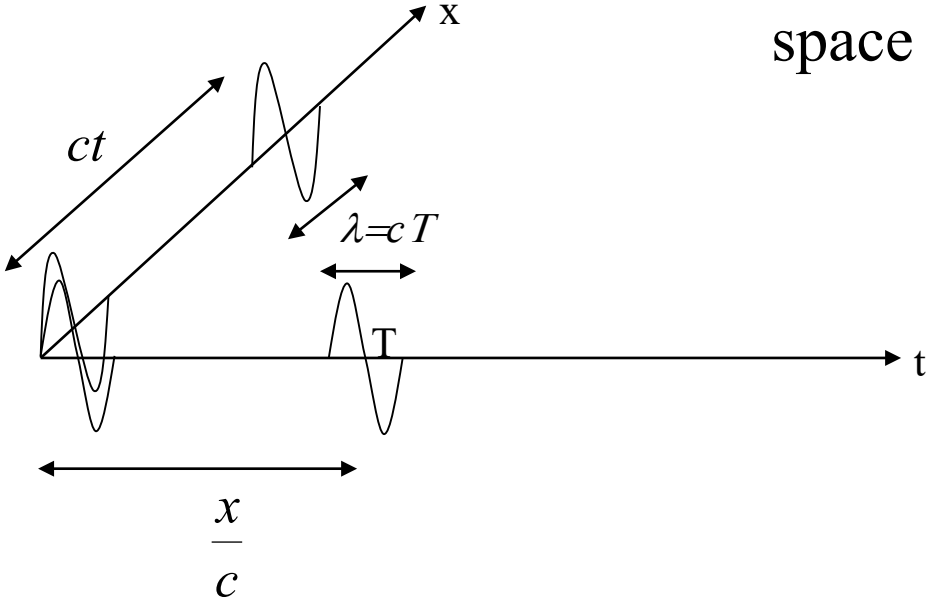
$$\varphi = \omega t - kx$$

Phase

Correspondance

Time

space

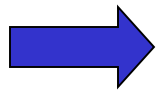


Complex notation

$$u_C = Ae^{j(\omega t - \phi)}$$

$$v_C = Be^{j(\omega t - \psi)}$$

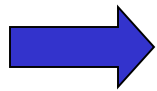
$$u = \Re(u_C)$$



derivative

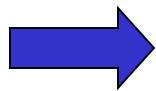
$$\frac{du_C}{dt} = j\omega u_C$$

$$\frac{du_C}{dx} = -jku_C$$



modulus

$$A^2 = u_C u_C^*$$

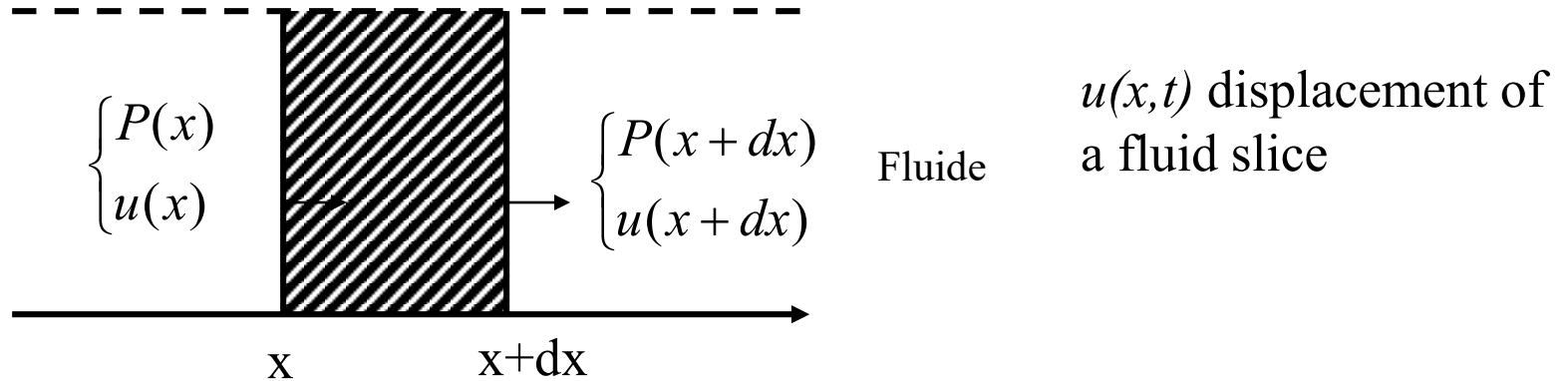


Mean value

$$\langle u(t)v(t) \rangle = \frac{1}{2}(u_C \cdot v_C^*)$$

Exercise : demonstrate the relation above

One dimensional equations of acoustic



Dynamic equation

$$dx\rho_0 \frac{\partial^2 u}{\partial t^2} = P(x) - P(x+dx)$$

$$P(x) - P(x+dx) = -dP = -\frac{\partial P}{\partial x} dx \quad \Rightarrow$$

$$\boxed{\rho_0 \frac{\partial^2 u}{\partial t^2} = -\frac{\partial P}{\partial x}}$$


Relation between P and u ?

Thermo says : the compressibility is

$$\chi = -\frac{1}{V} \frac{\partial V}{\partial P}$$

Pressure excess

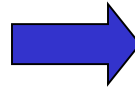
hence $\partial P = -\frac{1}{\chi} \frac{\partial V}{V}$



If A is lateral surface

$$\partial V = A[u(x+dx) - u(x)]$$

et $V = A dx$

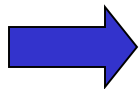


$$\frac{\partial V}{V} = \frac{\partial u}{\partial x}$$

Local displacement

=>

$$\partial P = -\frac{1}{\chi} \frac{\partial u}{\partial x}$$



Wave equation

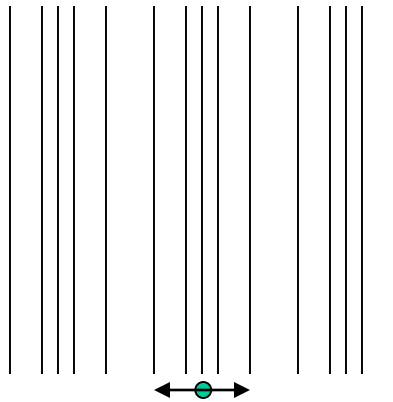
$$\frac{\partial^2 u}{\partial t^2} = \frac{1}{\rho_0 \chi} \frac{\partial^2 u}{\partial x^2} = c^2 \frac{\partial^2 u}{\partial x^2}$$

with

$$c = \sqrt{\frac{1}{\rho_0 \chi}}$$

Celerity of the acoustic wave

Fundamental laws wave propagation in a perfect, non turbulent fluid



P, \dot{v}, ρ

P , pressure

\dot{v} Particle velocity

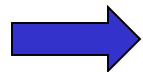
ρ , density

F an external force and q an acoustic source

Relations between these 3 variables and the sources ?

Hypothesis :

- small acoustic perturbations
- Adiabatic phenomenon
-



3 unknown



3 équations

Linear acoustic equations

$$P = c^2 \rho$$

$$\rho_0 \frac{d\vec{v}}{dt} = -\text{grad}(P) + \rho_0 \vec{F}$$

$$\frac{\partial \rho}{\partial t} + \rho_0 \text{div}(\vec{v}) = \rho_0 q$$

State (1)

Euler (2)

Mass (3)

Rq: $\dot{\vec{v}} = -\text{grad}(\phi) + \text{rot}(\psi)$

Vortex potential ($\text{div } \vec{v} = 0$)

Velocity potential ($\text{rot } \vec{v} = 0$)

Very often $\dot{\vec{v}} = -\text{grad}(\phi)$

No vortex

Wave equation

$$\begin{cases} \Delta P - \frac{1}{c^2} \frac{\partial^2 P}{\partial t^2} = \rho_0 (\operatorname{div}(\vec{F})) - \frac{\partial q}{\partial t} \\ \Delta \vec{v} + \operatorname{rot} \operatorname{rot} \vec{v} - \frac{1}{c^2} \frac{\partial^2 \vec{v}}{\partial t^2} = \operatorname{grad} q - \frac{1}{c^2} \frac{\partial \vec{F}}{\partial t} \end{cases}$$

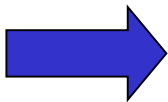
Exercice : with eq 1,2,3 show that we obtain these wave equations

□ □ □ □ Very complicated

if

$$\begin{aligned} \vec{v} &= -\operatorname{grad}(\phi) \\ \vec{F} &= -\operatorname{grad}(U) \end{aligned} \quad \text{then}$$

$$\begin{aligned} \Delta \phi - \frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} &= -q + \frac{\partial U}{\partial t} \\ P &= \rho_0 \frac{\partial \phi}{\partial t} - \rho_0 U \quad \text{Euler} \end{aligned}$$



ϕ gives v et P

Without sources

$$\left| \begin{array}{l} \Delta\phi - \frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} = 0 \\ \Delta P - \frac{1}{c^2} \frac{\partial^2 P}{\partial t^2} = 0 \\ \Delta \vec{v} - \frac{1}{c^2} \frac{\partial^2 \vec{v}}{\partial t^2} = 0 \end{array} \right.$$

$$\rho_0 \frac{d\vec{v}}{dt} = -\text{grad}(P)$$

General form

$$\phi(x, t) = F\left(t \pm \frac{x}{c}\right) = f((x \pm ct))$$

Exo : demonstrate that the above form is solution

For a plane wave propagating in the x direction

$$\begin{aligned} v_x &= V_M e^{j(\omega t - kx)} && \text{Toward } x > 0 \\ j\omega\rho_0 v_x &= -\frac{\partial P}{\partial x} = jkP = j\frac{\omega}{c}P \\ P &= \rho_0 c v_x \end{aligned}$$

Energy considerations

Kinetic energy density $E_C = \frac{1}{2} \rho_0 v^2$ [J/m³]

Potential energy $dW = -PdV$ [J]

Remember the state equation

$$P = c^2 \rho$$

 Per vol unit $dE_{P=} \frac{dW}{V} = P \frac{d\rho}{\rho_0} = \frac{c^2}{\rho_0} \rho d\rho$

then $E_P = \frac{1}{2} \frac{c^2}{\rho_0} \rho^2 = \frac{1}{2} \frac{P^2}{\rho_0 c^2}$

Energy density (J/m³) :

$$E_I = E_C + E_{P=} = \frac{1}{2} \rho_0 v^2 + \frac{1}{2} \frac{P^2}{\rho_0 c^2} \quad \text{In J/m}^3$$

Acoustic intensity:

$$v = V_m \sin(\omega t - kx) \quad P = P_m \sin(\omega t - kx)$$

The acoustic intensity is $I = Pv$ *in W/m^2*

$$I = P_m V_m \sin^2(\omega t - kx)$$

$$\langle I \rangle = \frac{P_m V_m}{2}$$

Harmonic regime

$$\tilde{v} = V_m e^{i(\omega t - kx)} \quad \tilde{P} = P_m e^{i(\omega t - kx)}$$

Complex quantities contain information on the magnitudes and phases of the wave.

Average Intensity

$$\langle I \rangle = \frac{\tilde{P} \tilde{v}^*}{2}$$

Rq 1: acoustic impedance

$$Z = \frac{P}{v} = |Z|e^{j\phi}$$

For a plane wave

$$Z = \rho_0 c$$

Acoustic impedance in $\text{kg m}^{-2} \text{s}^{-1}$
Or Rayl (Lord Rayleigh)

Rq 2:
$$E_I = E_C + E_{P=} \frac{1}{2} \rho_0 v^2 + \frac{1}{2} \frac{P^2}{\rho_0 c^2}$$

$$E_I = \frac{1}{2} \rho_0 v^2 + \frac{1}{2} \frac{P^2}{\rho_0 c^2} = \frac{1}{2c} \frac{P^2}{Z} + \frac{1}{2c} \frac{P^2}{Z} = \frac{P^2}{cZ}$$

$$\langle I \rangle = \frac{Pv^*}{2} = \frac{1}{2} Z v^2 = \frac{1}{2} \frac{P^2}{Z} = cE_I$$

$$\text{RMS pressure} \quad P_{rms} = \frac{P_m}{\sqrt{2}} \qquad \text{RMS velocity} \quad v_{rms} = \frac{v_m}{\sqrt{2}}$$

$$E_I = \rho_0 v_{rms}^2 + \frac{P_{rms}^2}{\rho_0 c^2}$$

P and v are conjugated variables like U et I in electricity

Electricity

$$U = RI$$

$$P = UI$$

$$P = \frac{1}{2} RI^2$$

$$P = \frac{1}{2} \frac{U^2}{R}$$

Acoustic

$$P = Zv$$

$$I = Pv$$

$$I = \frac{1}{2} Zv^2$$

$$I = \frac{1}{2} \frac{P^2}{Z}$$

sinus

Order of magnitude

Audible sounds Air: $\rho = 1,28 \text{ kg/m}^3$ $c = 340 \text{ m/s}$
 $Z = 440 \text{ Rayl}$

Limit of sensitivity of the human ear 10^{-12} W/m^2 or $P_r = 2 \cdot 10^{-5} \text{ Pa}$ à 1 kHz

$$I_{dB} = 10 \log_{10} \frac{I}{I_s} \quad \text{ou} \quad I_{SPL} = 20 \log_{10} \frac{P_{eff}}{P_r}$$

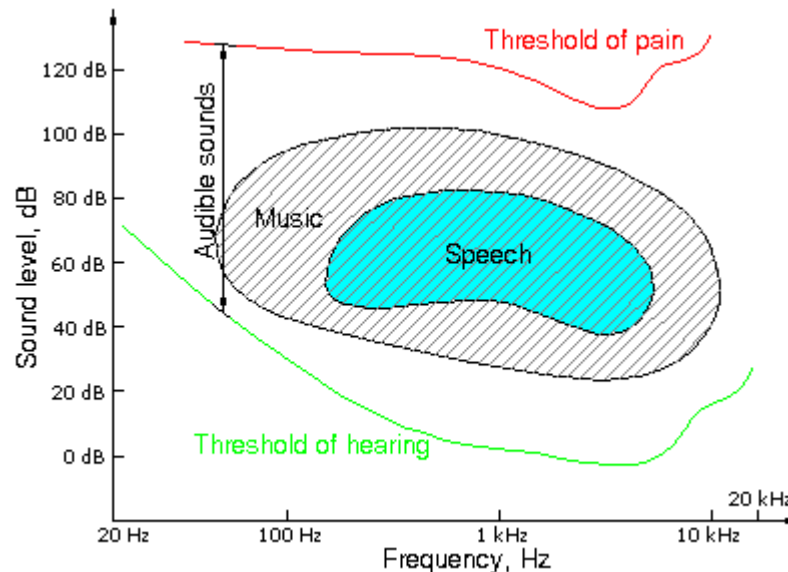
Quite house

air sound 40 à 50 dB

Conversation : 60 dB soit 10^{-6} W/m^2

concert: 80 à 90 dB

aircraft: 120 dB threshold of pain



Exercice : Calculate the mean displacement and velocity
 $f = 1 \text{ kHz}$ $I = 60 \text{ dB}$, in air

$$\text{In } W/m^2 \quad \langle I \rangle = 10^{-6} / 10^{-12}$$

The complex particle velocity is $v = V_m e^{i(\omega t - kx)}$,

In air $Z = \rho_0 c = 440 \text{ Rayl}$ and thus the intensity is :

$$\langle I \rangle = \frac{1}{2} Z v^2 = \frac{1}{2} Z V_m^2$$

The particle velocity is

$$V_m = \sqrt{\frac{2\langle I \rangle}{Z}} = \sqrt{\frac{2 * 10^{-6}}{440}} = 67 \mu m/s$$

$$U_m = \frac{V_m}{\omega} = \frac{V_m}{2\pi f} = 10 \text{ nm!!!}$$

In water Calculate the mean displacement and velocity

$f = 5 \text{ MHz}$ $I = I = I_{\max} = 100 \text{ mW/cm}^2$ in water

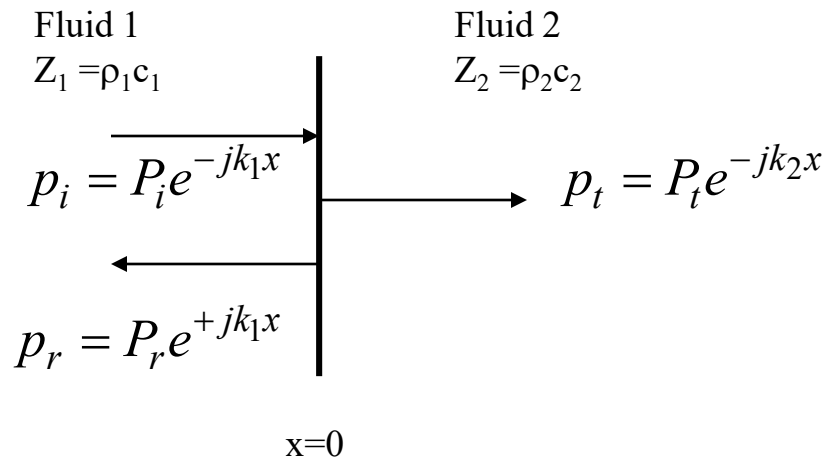
Density $\rho = 1000 \text{ kg/m}^3$ Celerity $c = 1500 \text{ m/s}$

$$V_m = 3.6 \text{ cm/s}$$

$$U_m = 1.2 \text{ nm}$$

Transmission and reflexion (plane waves)

Normal incidence



Reflexion coefficient

$$r = \frac{p_r}{p_i}$$

Transmission Coefficient

$$t = \frac{p_t}{p_i}$$

En $x = 0$: continuity equations

$$p_i + p_r = p_t$$

$$v_i + v_r = v_t$$

With $p = \pm Zv$

Depending on the direction

$$\frac{p_i}{Z_1} - \frac{p_r}{Z_1} = \frac{p_t}{Z_2}$$

Amplitude transmission coefficient

$$r = \frac{Z_2 - Z_1}{Z_2 + Z_1} \quad t = \frac{2Z_2}{Z_2 + Z_1}$$

Energy transmission coefficient

$$R = \frac{I_r}{I_i} = \frac{p_r^2 Z_1}{Z_1 p_i^2} = r^2 \quad R = \left(\frac{Z_2 - Z_1}{Z_2 + Z_1} \right)^2$$
$$T = \frac{I_t}{I_i} = \frac{p_t^2 Z_1}{Z_2 p_i^2} = \frac{Z_1}{Z_2} t^2 \quad T = \frac{4Z_1 Z_2}{(Z_2 + Z_1)^2}$$

with: $R+T=1$ energy is conservative

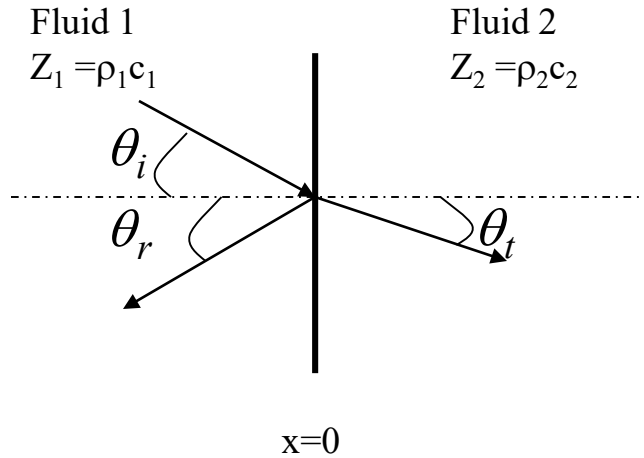
Exercice: Determine the transmission coefficient between air and water

In air $Z = 440 \text{ Rayl}$ in water $Z = 1.5 \cdot 10^6 \text{ Rayl}$ thus

$R = 0.9988$ and $T = 0.0012$

In dB, the attenuation due to transmission is close to 30 dB

Oblique incidence



Like in optics

$$\theta_i = \theta_r$$

and

$$\frac{\sin \theta_i}{c_1} = \frac{\sin \theta_t}{c_2}$$

if $\sin \theta_i = \frac{c_1}{c_2}$

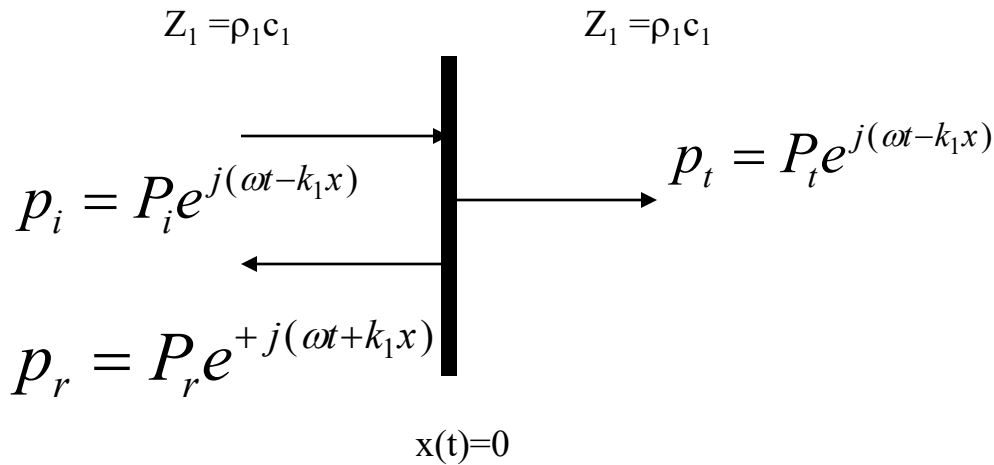
Total reflexion

$$r = \frac{Z_2 \cos \theta_i - Z_1 \cos \theta_t}{Z_2 \cos \theta_i + Z_1 \cos \theta_t}$$

$$t = \frac{2Z_2 \cos \theta_i}{Z_2 \cos \theta_i + Z_1 \cos \theta_t}$$

Exemple : propagation through a rigid wall surface density μ in kg/m^2

Normal incidence



$$v_i + v_r = v_t \quad \frac{p_i}{Z_1} - \frac{p_r}{Z_1} = \frac{p_t}{Z_2} = v(t)$$

thus
$$p_r = p_i - \frac{Z_1 p_t}{Z_2}$$

Dynamic equation

$$p_i + p_r - p_t = \mu \frac{dv(t)}{dt}$$

If $Z_1 = Z_2$

For a sinusoidal wave

$$p_i + p_i - p_t - p_t = \mu \frac{dv(t)}{dt} = j\omega\mu \frac{p_t}{Z_1}$$

$$p_i - p_t = j \frac{\omega\mu}{2} \frac{p_t}{Z_1}$$

$$t\left(1 + j\frac{\omega\mu}{2} \frac{1}{Z_1}\right) = 1$$

$$t = \frac{p_t}{p_r} = \frac{1}{1 + j\frac{\omega\mu}{2\rho_1 c_1}}$$

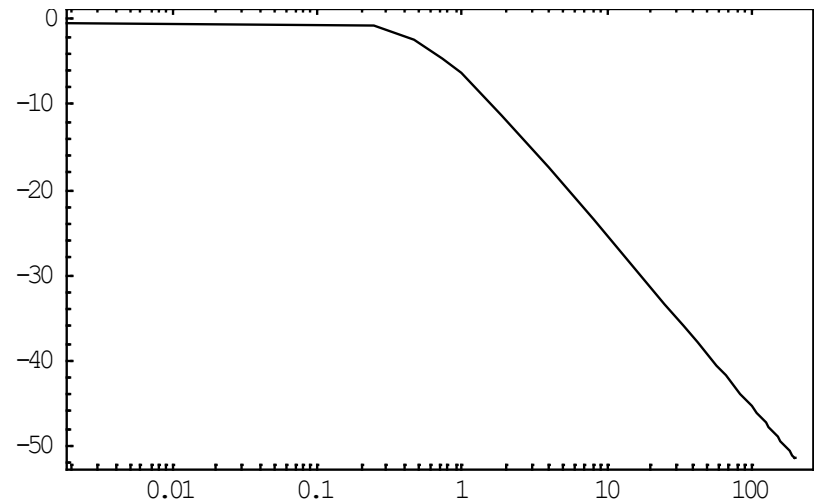
Low pass filter

$$\text{Et } T = \left|\frac{p_t}{p_r}\right|^2 = \frac{1}{1 + \left(\frac{\omega\mu}{2\rho_1 c_1}\right)^2}$$

Numerical application

Concrete wall: $e = 10 \text{ cm}$, $\rho = 2600 \text{ kg/m}^3 \Rightarrow \mu = 260 \text{ kg/m}^2$

$T = 46 \text{ dB}$ à 100 Hz



Attenuation due to viscosity. η, μ

Dynamic equation

$$\rho_0 \frac{d\vec{v}}{dt} = -\text{grad}(P) + \left(\eta + \frac{4}{3}\mu\right)\Delta\vec{v} \quad \text{i.e.} \quad \frac{d^2\vec{u}}{dt^2} = c^2\Delta\vec{u} + \frac{1}{\rho_0}\left(\eta + \frac{4\mu}{3}\right)\Delta\frac{d\vec{u}}{dt}$$

Harmonic regime $u_x = U_M e^{j(\omega t - kx)}$

$$\frac{d^2\vec{u}}{dt^2} = \left(c^2 + j\omega \frac{1}{\rho_0}\left(\eta + \frac{4\mu}{3}\right)\right)\Delta\vec{u}$$

Complex celerity

$$c^{2*} = c^2 \left(1 + j\omega \frac{1}{\rho_0 c^2} \left(\eta + \frac{4\mu}{3}\right)\right) \quad \text{First order dev. : } c = c \left(1 + j\omega \frac{1}{2\rho_0 c^2} \left(\eta + \frac{4\mu}{3}\right)\right)$$

Complex wave number

$$k^* = \frac{\omega}{c^*} = \frac{\omega}{c} \left(1 - j\omega \frac{1}{2\rho_0 c^2} \left(\eta + \frac{4\mu}{3}\right)\right)$$

Expression of the particule displacement

$$u_x = U_M e^{j(\omega t - k^* x)} \quad \text{Soit} \quad u_x = U_M e^{j(\omega t - \frac{\omega}{c} x)} e^{-\alpha x}$$

Attenuation coefficient
Propagation
Attenuation

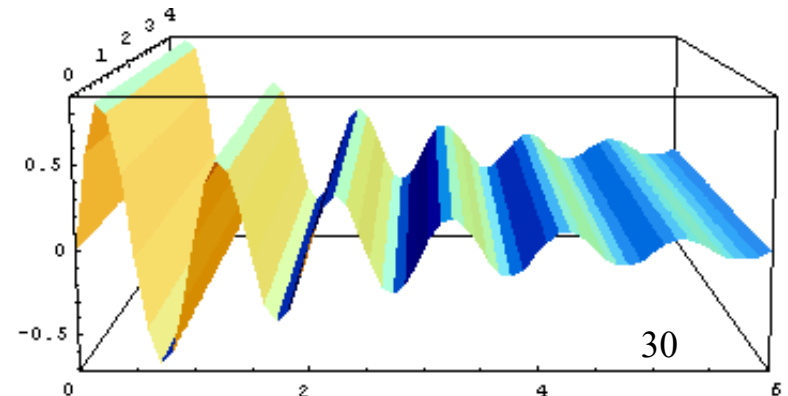
$$\alpha = \frac{\omega^2}{2\rho_0 c^3} \left(\eta + \frac{4\mu}{3} \right) x \quad \text{Neper/m} \quad \text{In dB}$$

$$A(\text{dB}) = 20 \log_{10}(e^{-\alpha x}) = 8.686 \alpha x$$

$$A(\text{dB}) = 8,686 \frac{\omega^2}{2\rho_0 c^3} \left(\eta + \frac{4\mu}{3} \right) x$$

Attenuation proportionnal to the square frequency

Other causes : thermal conductivity
 Molecular relaxation



Attenuation in air in dB/m at 1 kHz, 1 MHz,
 $\mu = 1.910 \cdot 10^{-5}$, $\eta = 0.6\mu$

What is the viscous interaction length?

$$\alpha = \frac{\omega^2}{2\rho_0 c^3} \left(\eta + \frac{4\mu}{3} \right) x$$

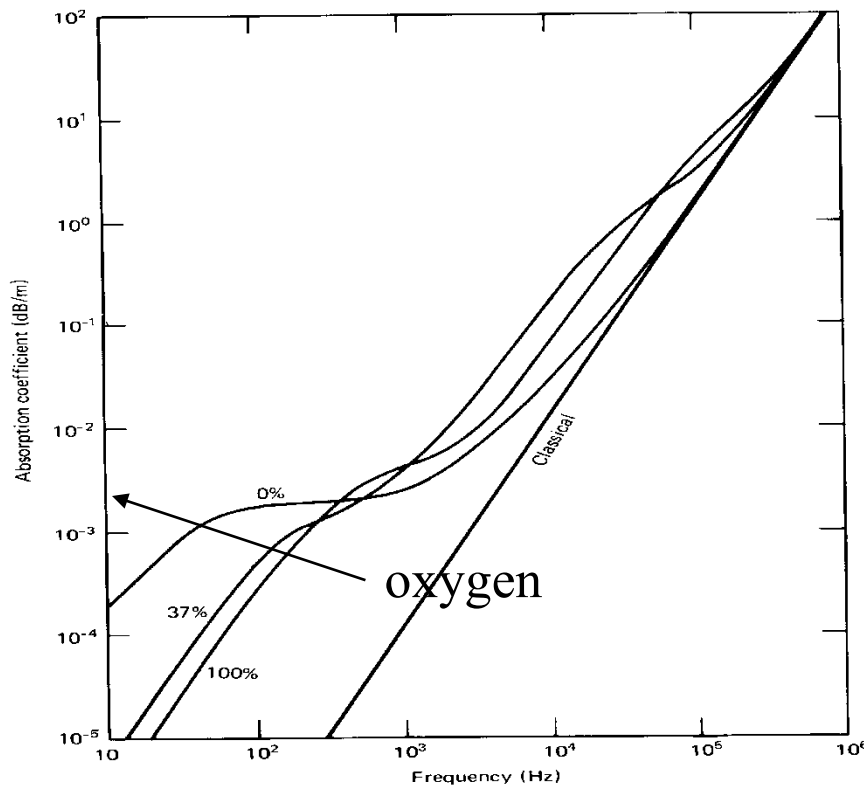


Fig. 7.3. Absorption of sound in air at 20°C as a function of frequency. (After Bass, et al.)

@ 1 kHz
 $\delta = 150 \text{ km!!}$
 @ 1 MHz
 $\delta = 0.15 \text{ m}$

Attenuation in liquids

$$\eta \ll \mu$$

$$\alpha \approx \frac{\omega^2}{\rho_0 c^3} \frac{2\mu}{3} x \quad \text{In Neper}$$

$$\alpha_{\text{water}} = 0.22 \text{ dB/m}$$

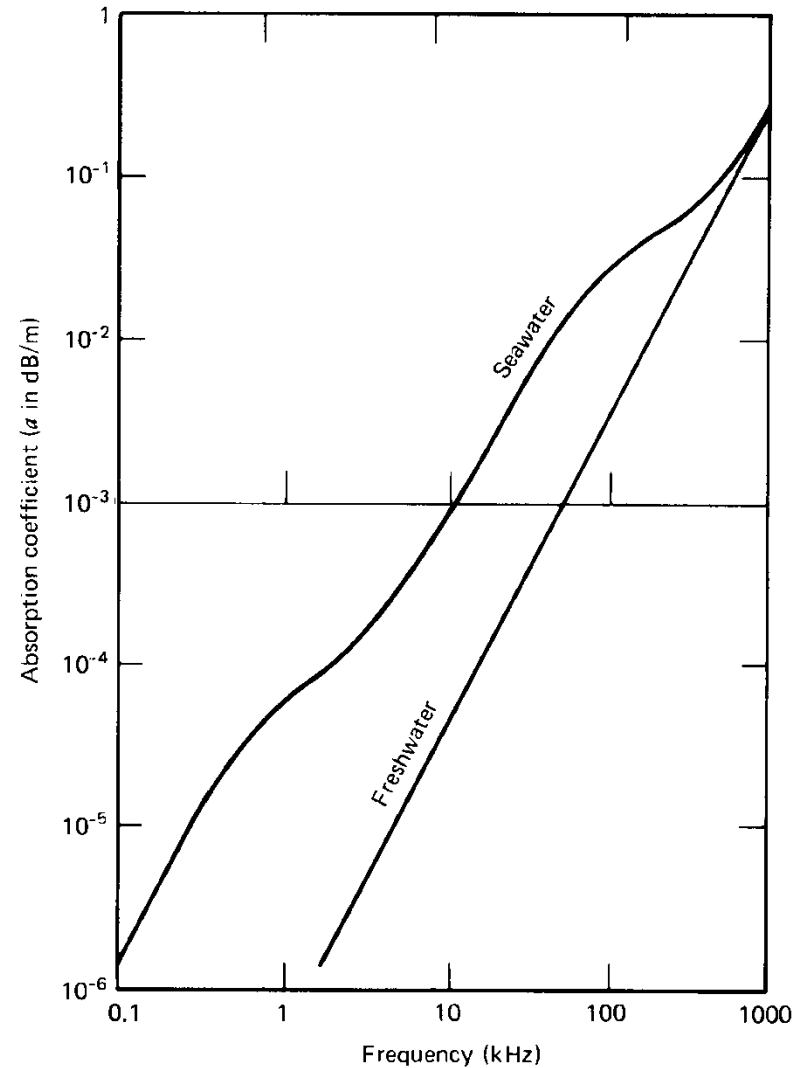


Fig. 7.5. Sound absorption in freshwater and in seawater (35 ppt salinity) at 5°C and 1 atm. (According to Fisher and Simmons.)

Conclusion

- Some concepts have been defined such as wave characteristics, acoustic impedance, transmission reflexion coefficients, attenuation factor...
- These concepts are essentials understand the interaction of a wave with mater, it is essential to know the essential characteristics of the wave celerity attenuation.
- We will use these concept to understand the electroacoustic behavior of a transducer and the wave interaction with materials

Few words on Piezoelectricity

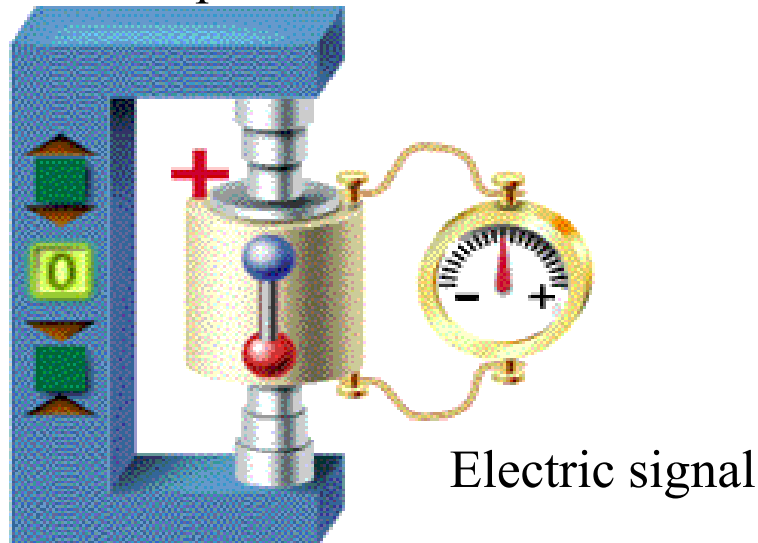
Définition

Material properties

Ferroelectric materials

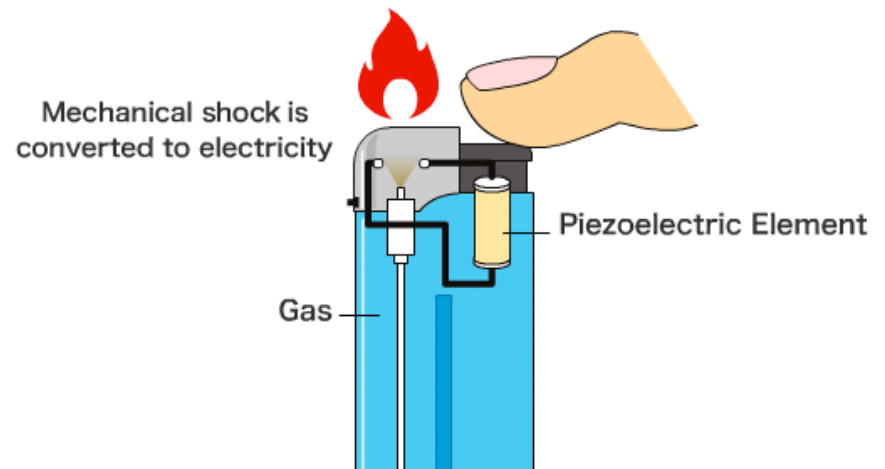
The piezoelectricity is the property of several materials that are able to generate an electric charge when mechanically deformed. Conversely, when an external electric field is applied to piezoelectric materials they mechanically deform.

Mecanic pressure



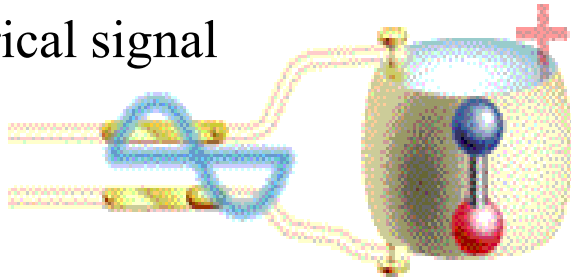
Direct effet

Lighter sensors



Mechanical to electrical energy conversion and conversely

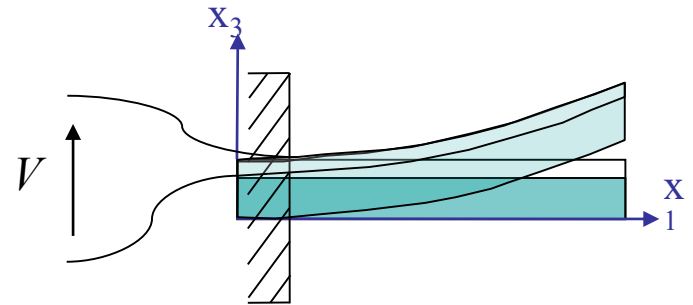
Electrical signal



Mechanical displacement

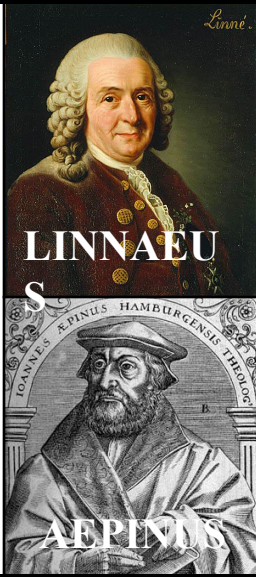
Inverse effect

actuators



Discovery of the piezoelectric effect

Pyroelectricity – studied c. 1750

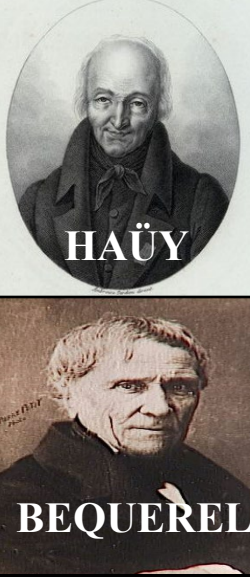


LINNAEUS

AEPINUS

The left panel contains two portraits. The top one is a color portrait of Carl Linnaeus, a Swedish naturalist, with the name 'LINNAEUS' printed below it. The bottom one is a black and white engraving of Johann Aepinus, a German physicist, with the name 'AEPINUS' printed below it. A green arrow points from this panel towards the right panel.

Piezoelectricity – early experimental work c. 1820



HAÜY

BEQUEREL

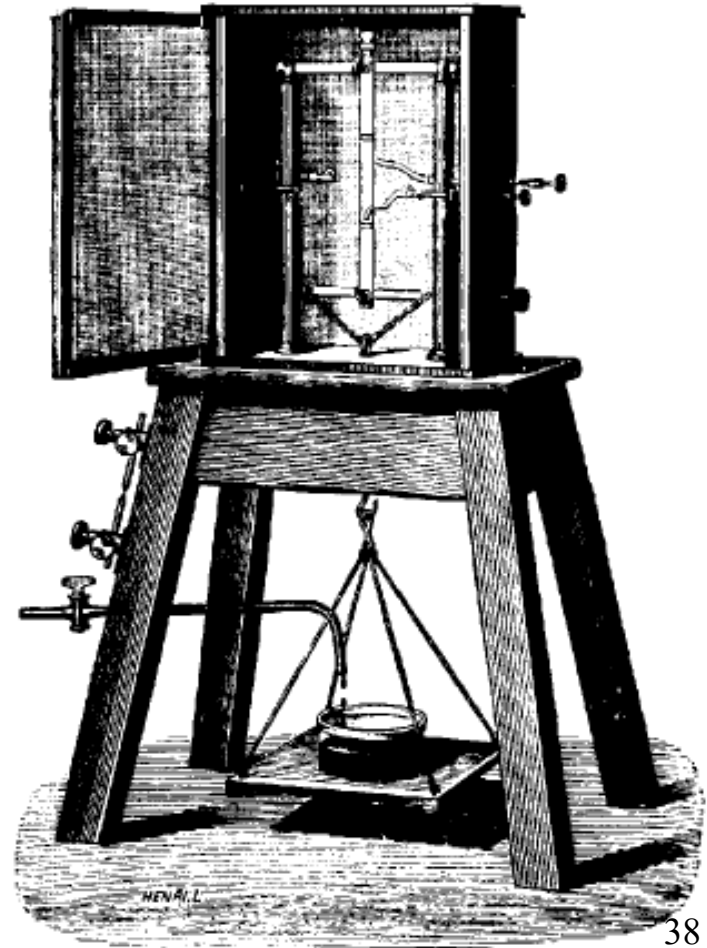
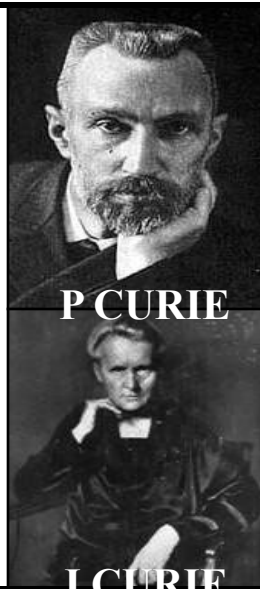
The right panel contains two portraits. The top one is a black and white oval portrait of Augustin Haüy, a French physicist, with the name 'HAÜY' printed below it. The bottom one is a color portrait of Antoine Becquerel, a French physicist, with the name 'BEQUEREL' printed below it.

Discovery of the piezoelectric effect

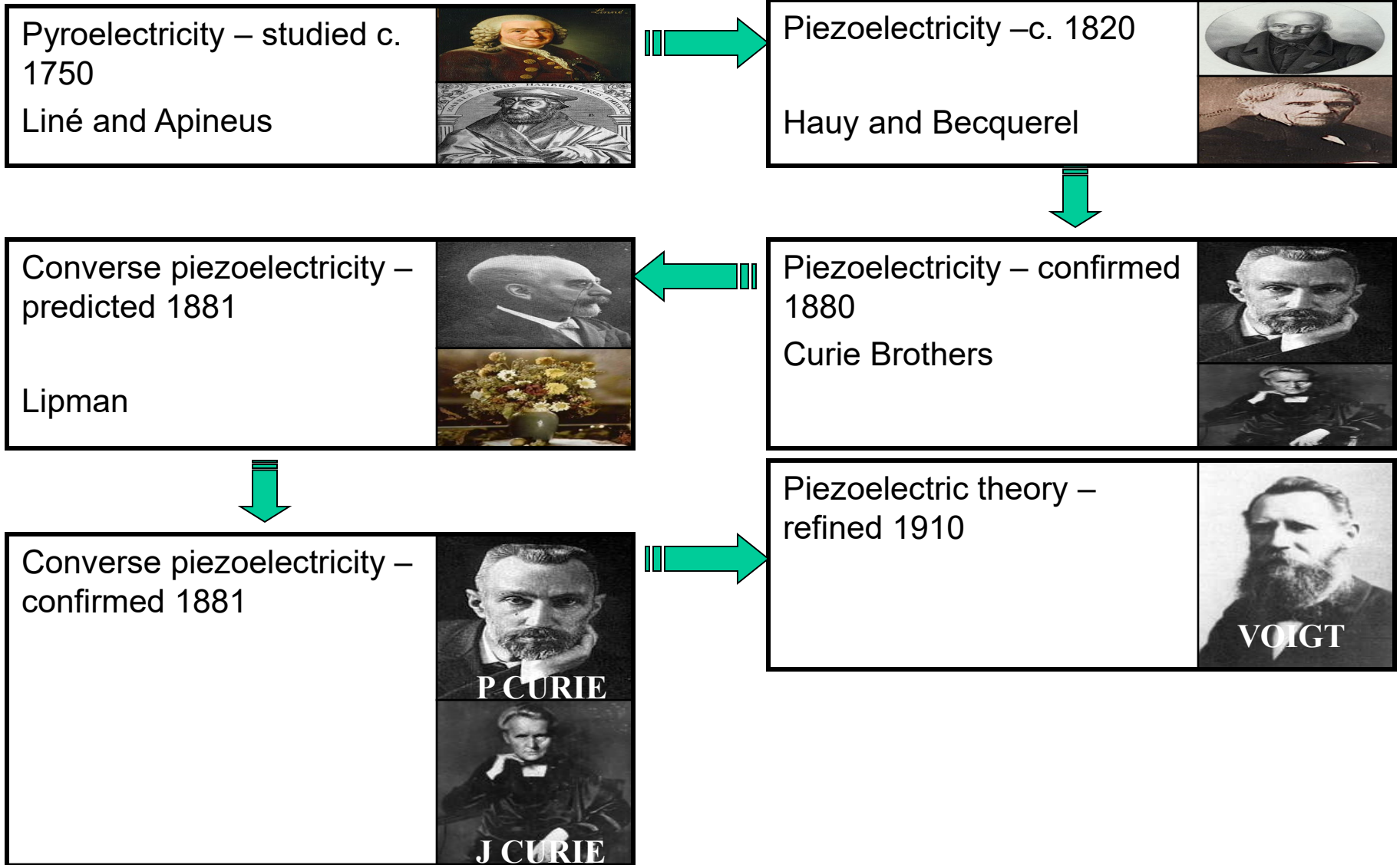
Curie brothers discover that several material are producing charges under mechanical deformation



Experimental evidence
in 1880



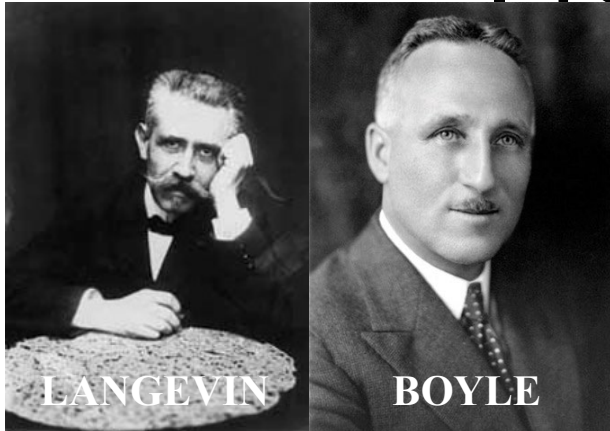
Discovery of the piezoelectric effect



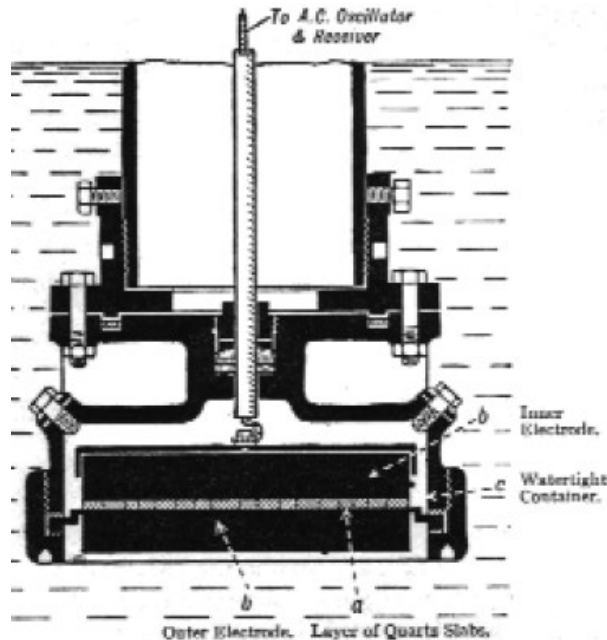
First World War

1916: ultrasonic transmission using
Quartz crystal

Work commissioned by the Anti-Submarine Division, in a joint research effort by the French, British and American lead to the development of early sonar (ASDIC)



6164_NoiseCollector_jillys_sonar.wav



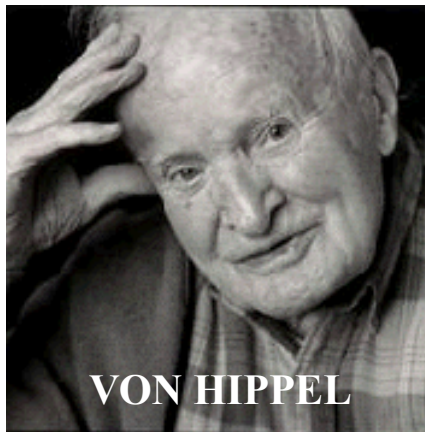
Quartz Transducer sandwich type h (A. B. Wood, A Textbook of Sound (1932))



6165_NoiseCollector_jillys_sonar⁴⁰2.wav

World War II & beyond

The origin of the high dielectric constant in barium titanate was discovered to be a result of its ferroelectric properties



XCVI. *Theory of Barium Titanate.*—Part I.

By A. F. DEVONSHIRE,

H. H. Wills Physical Laboratory, University of Bristol*.

[Received July 26, 1949.]

SUMMARY.

The theory of the dielectric and crystallographic properties of barium titanate is considered. By expanding the free energy as a function of polarization and strain and making reasonable assumptions about the coefficients, it is found possible to account for the various crystal transitions. Calculations are made of the dielectric constants, crystal strains, internal energy, and self polarization as functions of temperature. Finally relations are obtained between the coefficients in the free energy and the ionic force constants. These are used to estimate some of the coefficients which are not completely determined by experimental data.

High dielectric constant ceramics

A. von Hippel, R. G. Breckenridge,
F. G. Chesley, and Laszlo Tisza

LABORATORY FOR INSULATION RESEARCH, MASSACHUSETTS INSTITUTE OF TECHNOLOGY,
CAMBRIDGE, MASS.

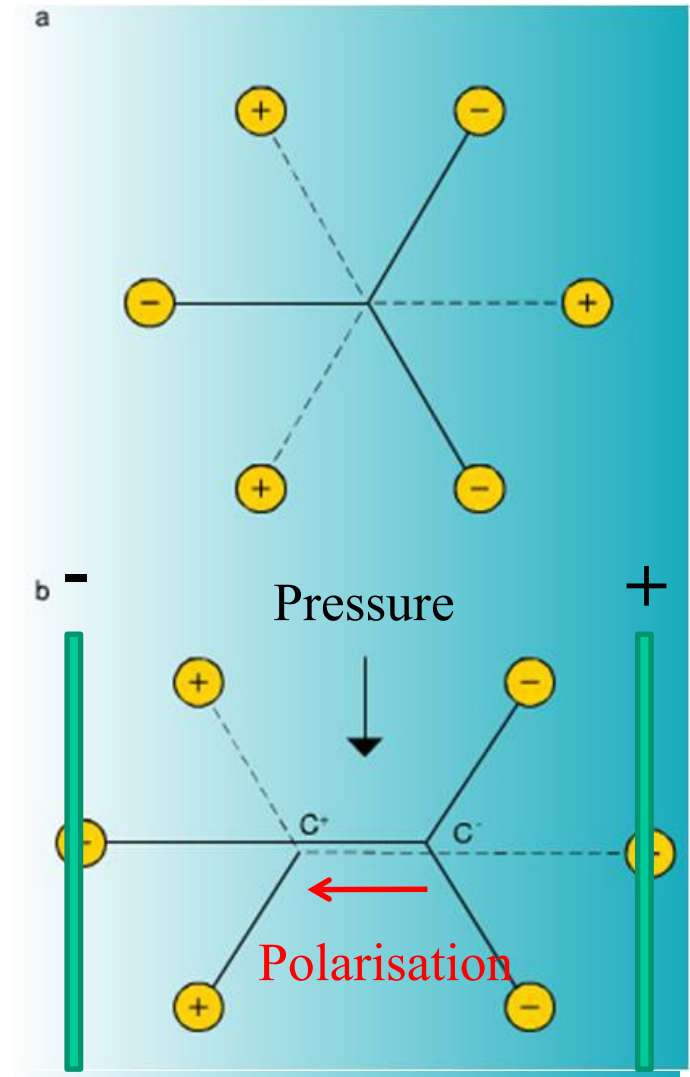
Results of an investigation are presented in this paper on titanium dioxide and the alkaline earth titanates, including some mixtures and solid solutions of the barium and strontium compounds. Dielectric measurements over a wide range of frequencies, temperatures, and voltages, and thermal expansion and x-ray studies have yielded a rather complete picture of the properties and usefulness of these titania ceramics. Barium titanate and the

barium-strontium titanate solid solutions prove to be a new class of ferroelectric materials. Their peculiar dielectric behavior was first noted by the Titanium Alloy Manufacturing Company, and this behavior proves to be connected with a lattice transition from pseudocubic to cubic. Additional maxima have been found in the dielectric characteristics which correspond to transitions of the second order. These maxima are being studied further.

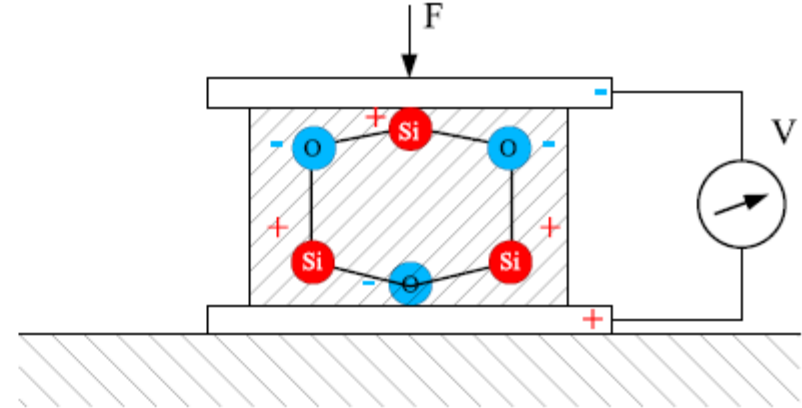
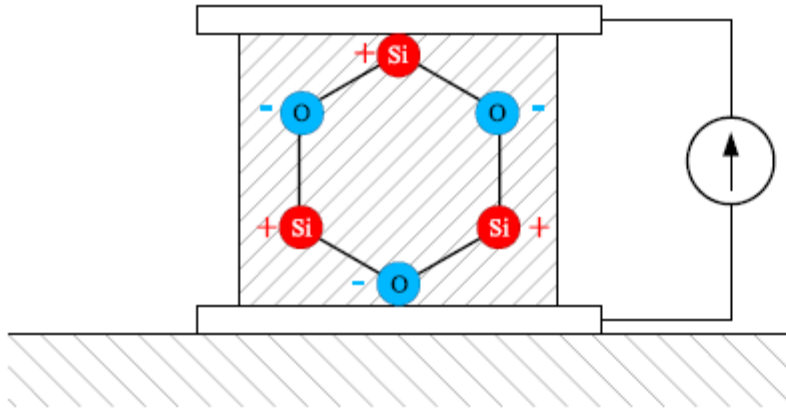
A complete theory predicting the structural phase transitions in BaTiO_3 was proposed by Devonshire

Why in several materials

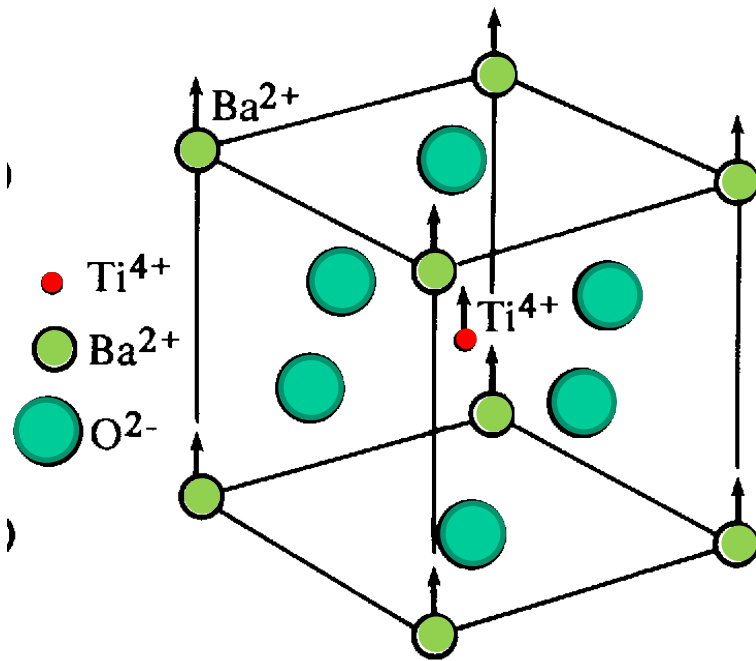
Piezoelectricity
appears in non centro
symmetrical crystals



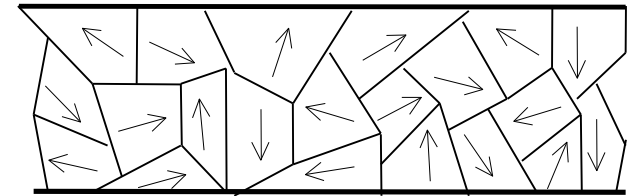
Example: Quartz crystal charges appear on top and bottom faces



The ferroelectricity and the Piezoelectricity barium titanate ceramics

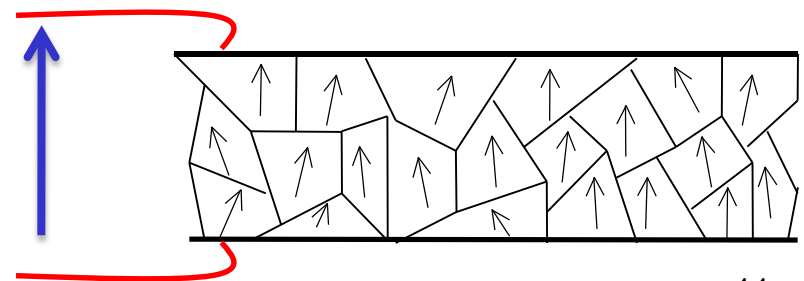


Before polarisation



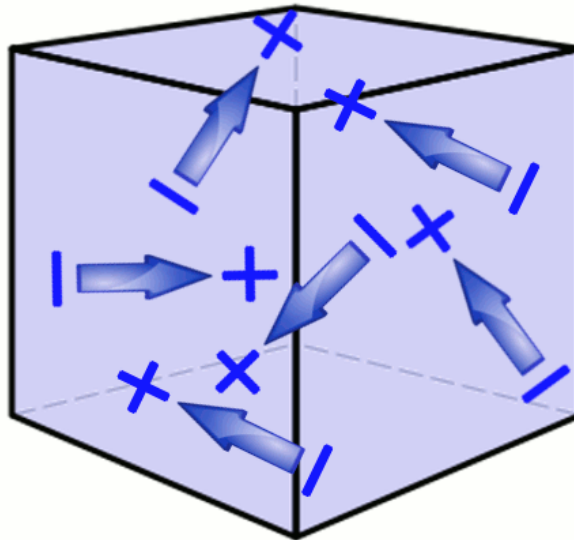
After polarisation

Voltage

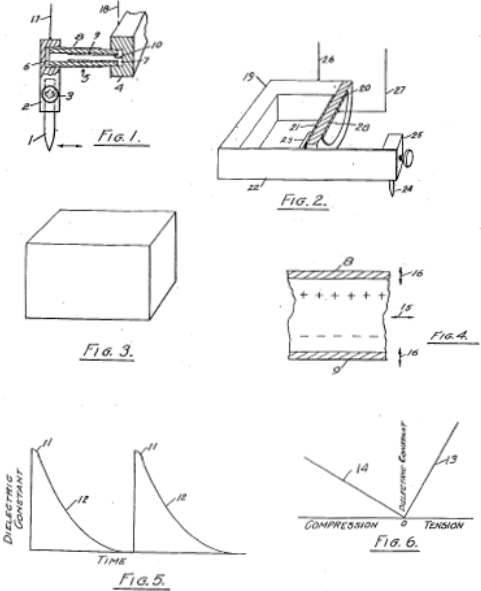


The development of ceramics

In 1949 a US patent was filed by Erie Resistor Corporation, for a piezo transducer that used the first poled ceramic samples.



Nov. 1, 1949. R. B. GRAY 2,486,560
 TRANSDUCER AND METHOD OF MAKING THE SAME
 Filed Sept. 20, 1946



UNITED STATES PATENT OFFICE

2,486,560

TRANSDUCER AND METHOD OF
 MAKING THE SAME

Robert B. Gray, Erie Pa., assignor to Erie Resistor
 Corporation, Erie, Pa., a corporation of Penn-
 sylvania

Application September 20, 1946, Serial No. 698,374

15 Claims. (Cl. 171-337)

INVENTOR
 BY Robert B. Gray
 Ralph Hamman
 Attorney

PZT and PLZT ceramics



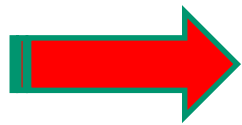
- Higher electromechanical coupling
- Higher TC values, resulting in higher operational and processing temperatures
- Easily poled for applications, versatility
- Possess a wide range of dielectric constants across the phase diagram



Hans Jaffe (far right)



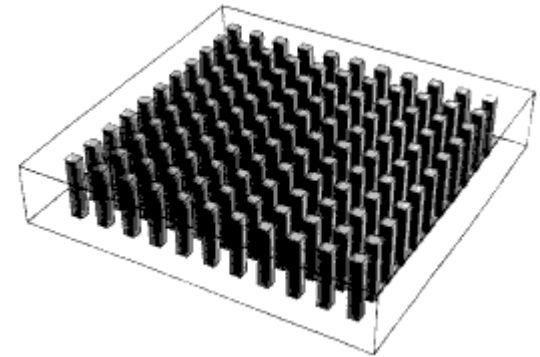
PZT ceramics
(www.fujicera.co.jp)



New technological processes Enhancement of the performances

1978 : composites materials

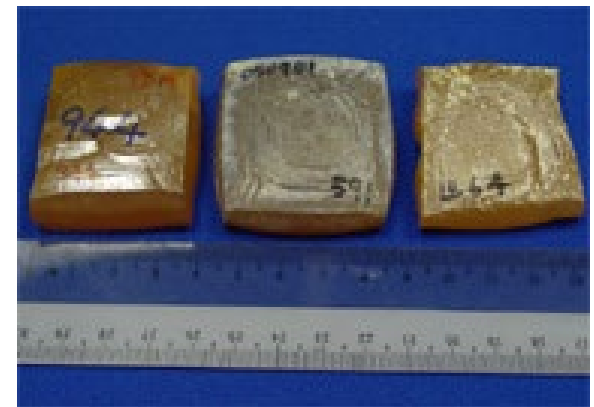
R.E. Newnham, D.P. Skinner, L.E. Cross,
Connectivity and piezoelectric-pyroelectric composites, *Materials Research Bulletin*, Volume 13, Issue 5, May 1978, Pages 525-536



<http://www.ndt.net/article/v06n08/fleury/fig1.gif>

90 Piezoelectric single crystals with giant properties

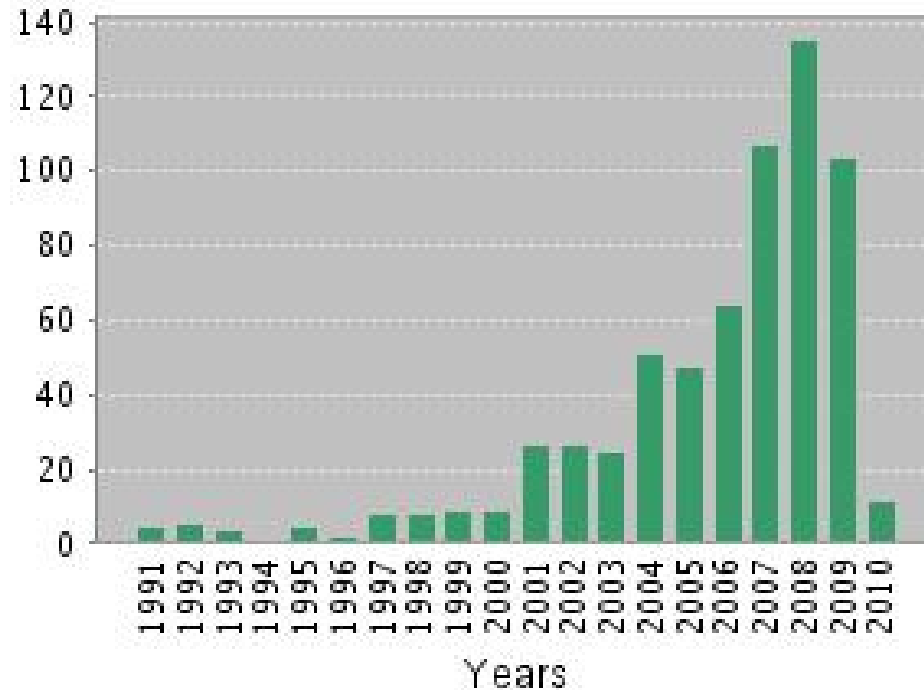
Jun Kuwata, Kenji Uchino, and Shoichiro Nomura,
Phase Transitions in the $\text{Pb}(\text{ZnInNb})\text{O}_3$ - PbTiO_3 System, *Ferroelectrics*, vol. 37, pp. 579-582, 1981
T. R. Shout, Z.P. Chang, N. Kim and S. Markgraf,
Dielectric Behavior of single Crystals Near the $(1-x)\text{Pb}(\text{Mg}_{1/3}\text{Nb}_{2/3})\text{O}_3$ $(x)\text{PbTiO}_3$ Morphotropic Phase Boundary, *Ferroelectr. Let.*, vol. 12, pp. 63-69, 1990.



<http://www.microfine-piezo.com/>

2000 : lead free ceramics

Published Items in Each Year

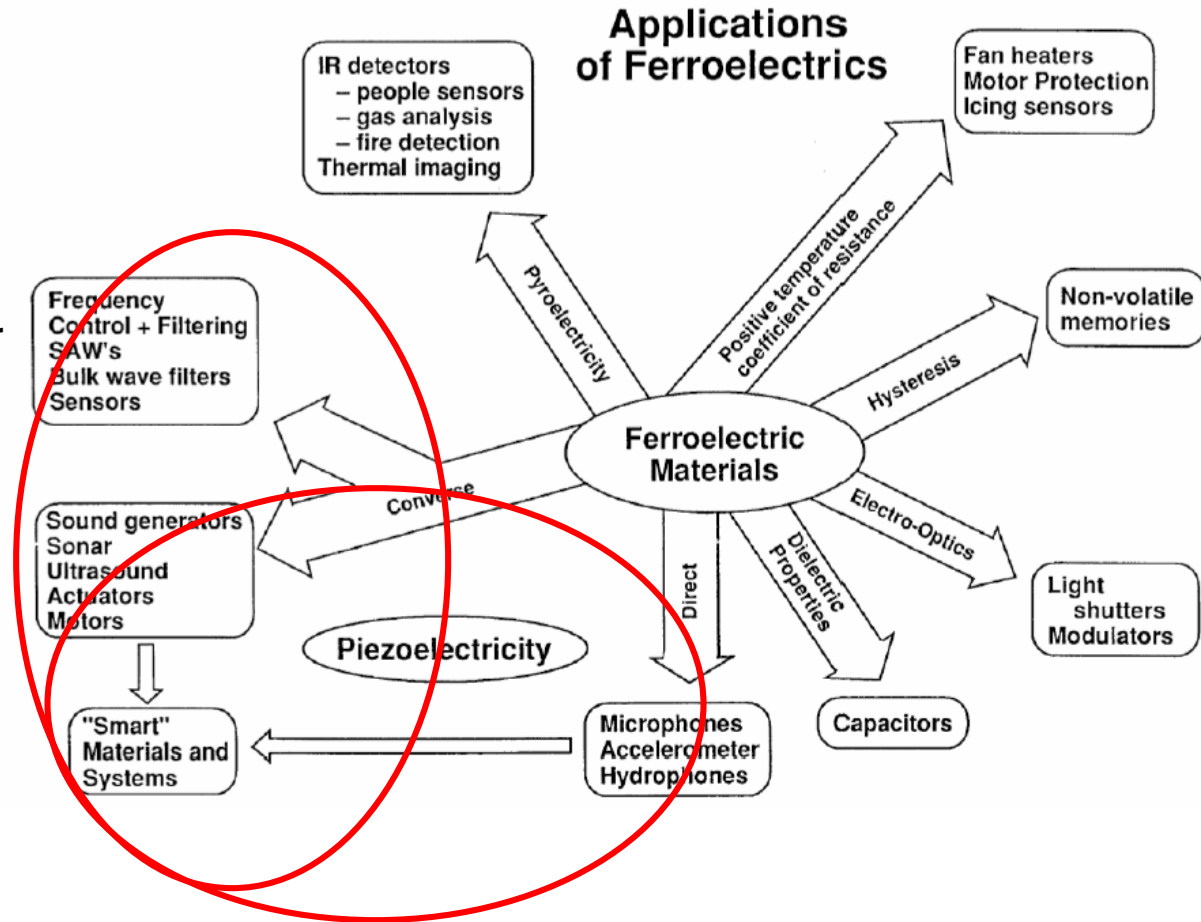


DIRECTIVE 2002/95/EC OF THE EUROPEAN PARLIAMENT AND OF THE COUNCIL
of 27 January 2003

on the restriction of the use of certain hazardous substances in electrical and electronic equipment

Accelerometers.

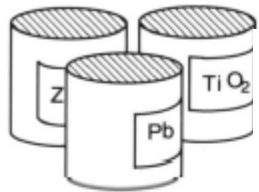
- Flow meters: Blood, industrial processes, waste water.
- Medical: Imaging, HIFU, IVUS, surgical knives, and cleaning of blood veins.
- Underwater acoustics: echosounders, sonar systems, fish-finders, seabed mapping.
- Industrial sensors based on ultrasound: Level control, detection, and identification.
- Hydrophones: Seismic, biologic, military, underwater communication.
- Inkjet printheads.
- Dental work: Removal of plaque.
- Alarm systems: Movement detectors, broken window sensors.
- NDT: Transducers for Non Destructive Testing.
- Musical instrument pickups.
- Acoustic emission transducers.
- Actuators.
- Micro positioning devices: Optics, scanning tunneling microscopes.
- Surface Acoustic Waves: Personal Computer touch screens, filters.



Material and properties

Fabrication process of ceramics

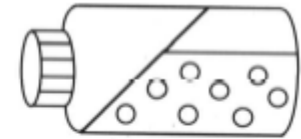
Powder



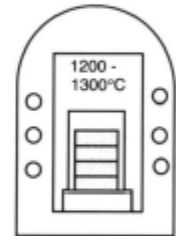
[1] RAW MATERIALS



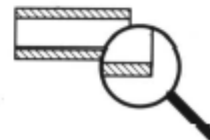
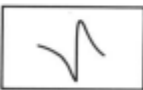
[5] G



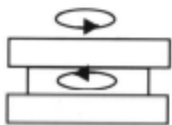
[4] MILLING



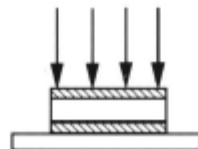
[8] SINTERING



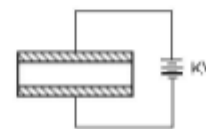
[12] FINAL INSPECTION



[9] GRINDING/POLISHING



[10] ELECTRODING



[11] POLING

Piezo ceramics

Shape:	Plate	Disc	Ring	Bar		
Polarization Direction						
Applied Field						
Voltage Output						
Mode of Vibration				(L) Length		
Displacement						
					Cylinder Wall Electrode	Plate Bender
Polarization Direction						
Applied Field						
Voltage Output						
Mode of Vibration						
Displacement						

Mains characteristics

- High stress
- Linear response
- Small deformations

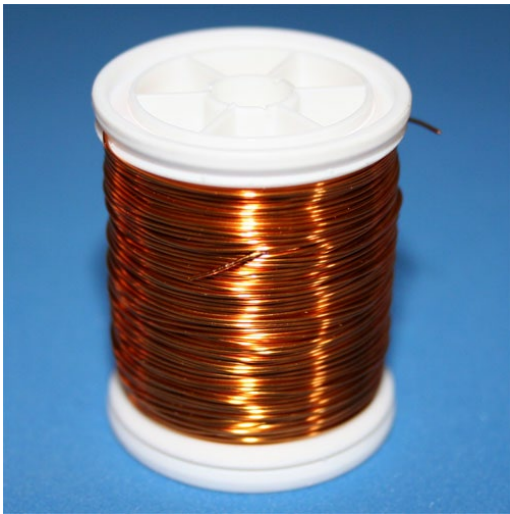
Piezo coefficient = 100 pC/N
 thickness = 1 mm
 Voltage = 10 V

displacement
 $\Delta l = 1 \text{ nm} !$

Some properties

Electrical

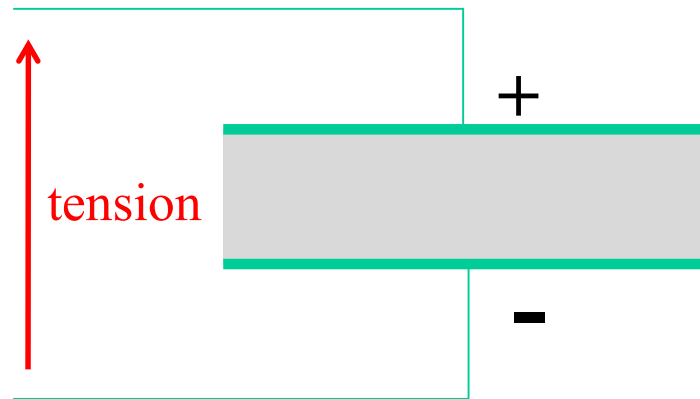
Piezoelectric material are insulators : characterised by their ability to accumulate charges



In a conducting metal electrons are free

In V/C

ϵ dielectric constant



In a insultator

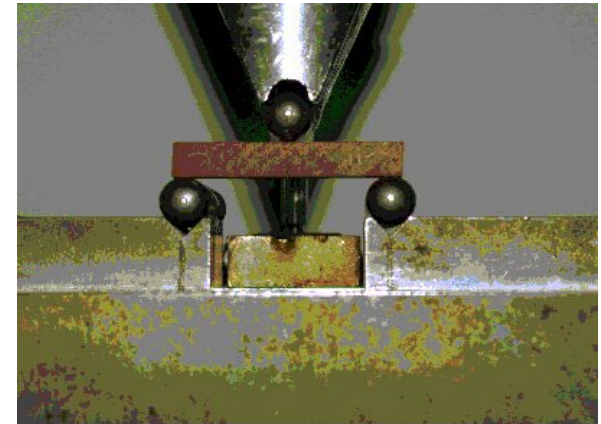
Electrons are located on electrodes

Mechanical Elasticity or stiffness

c coefficient



<http://www.insep.fr/FR/Sports/Gymnastique/PublishingImages/gymnastique.jpg>



<http://www.unilim.fr/theses/2003/sciences/2003limo0004/images/figurevi1.jpg>

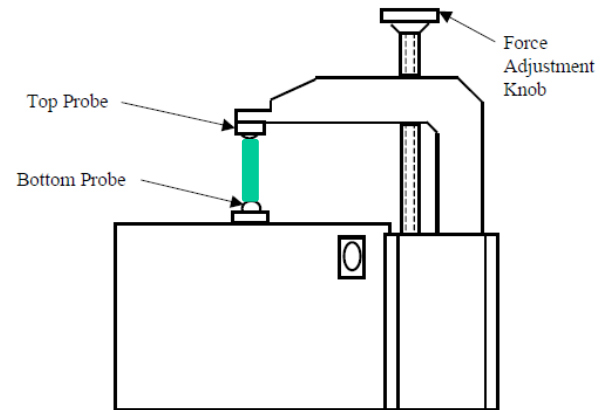
In Newton per m²

Electromechanical

Ability to accumulate charges for a given force



<http://www.piezotest.com/d33piezometer.htm>



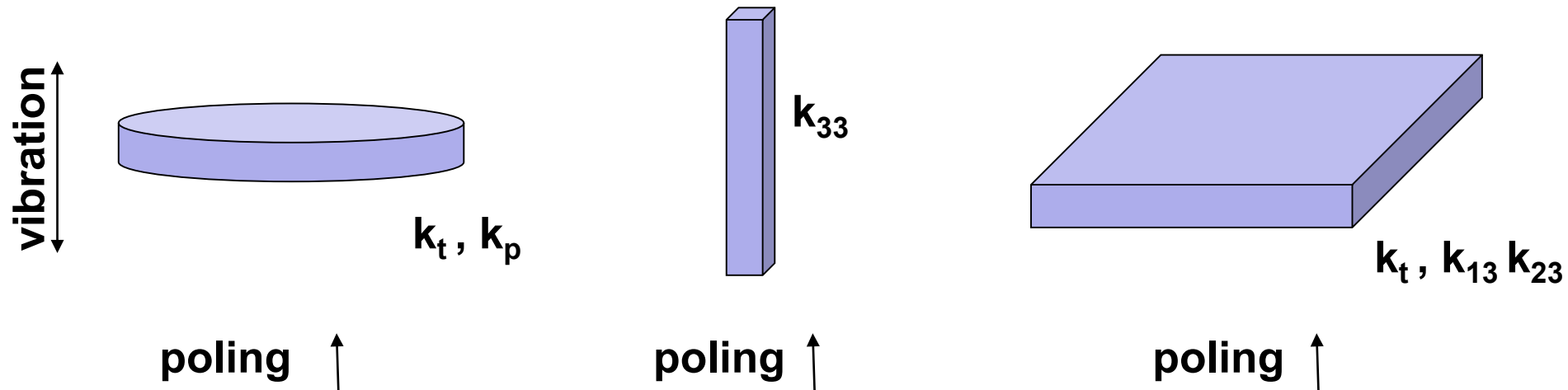
<http://www.kcftech.com/products/documents/PM3001.manual.pdf>

d coefficient

In Coulomb per Newton or m/Volt

Electromechanical coupling coefficient

The electromechanical coupling factor, k , of piezoelectric materials, determines the conversion efficiency of mechanical to electrical energy or electrical to mechanical energy.



The coupling factor, k , is associated with a vibration mode and thus do not have a unique expression, k is always bellow

Properties

Material	piezo. Coef. d_{33} (10^{-12} m/V)	Relative Permittivity	Stiffness coef. (10^9 N/m ²)	Density kg/m ³	Thicness coupling coef. (%)
Quartz [1]	2,3	4,5	80	2650	10
Barium Titanante (T57) [2]	105	700	112	5300	40
Lead Titanate (PZ34) [3]	50	210	140	7550	40
Lead Titanate Ziconate (PZ27) [3]	425	1800	113	7700	60
Piezo Polymers	26 [4]	5 [5]	10 [5]	2500 [5]	25 [5]

- [1] Morgan Matroc
- [2] Quartz et Silice
- [3] Ferroperm
- [4] Sasaki **ISSN** 0386-2186
- [5] Feuillard, thèse de doctorat

Conclusions

- Piezoelectric materials are used as sensor,actuators or transducers in many applications .
- Covered field : **microsystem, microelectronics medicine nuclear or aeronautics.**
- Very often hidden but essential in many applications**

- Piezo materials are known ...
- We focused mainly in transducer but they are used in many application
- Now it is time to see how we can use them to make transducers and to produce images...
- To be continued...

Thank you !